

ON THE GAUGE VERSION OF EXPONENTIAL MAP

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ABSTRACT. Let $P(M, G)$ be a principal bundle, Γ be a principal connection on P and Λ be a classical connection on M . For every $u \in P$, we construct a local map $\exp_u^{\Gamma, \Lambda}: T_u P \rightarrow P$. Then we clarify how these maps can be used for finding all natural operators transforming (Γ, Λ) into a principal connection of the r -th principal gauge prolongation $W^r P$ of P . We also deduce that $\exp^{\Gamma, \Lambda}$ coincides with the exponential map of a classical connection on P .
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INTRODUCTION

We study certain geometric properties of principal connections on higher order principal gauge prolongations of principal bundles that are related with the gauge theories of mathematical physics, [1, 3, 4, 5]. We generalize the concept of exponential map of a classical torsion-free connection Λ on a manifold M to the case of a principal G -bundle $p: P \rightarrow M$. We clarify that a principal connection Γ on P and Λ define a local map $\exp_u^{\Gamma, \Lambda}: T_u P \rightarrow P$ for every $u \in P$. Taking into account a linear frame v on M at $p(u)$, we modify $\exp_u^{\Gamma, \Lambda}$ into a local \mathcal{PB} -isomorphism $\exp_{v,u}^{\Gamma, \Lambda}: \mathbb{R}^m \times G \rightarrow P$, $m = \dim M$. The $(r+1)$ -jets of $\exp_{v,u}^{\Gamma, \Lambda}$ at $(0, e) \in \mathbb{R}^m \times G$ determine a reduction of the $(r+1)$ -st principal gauge prolongation $W^{r+1} P$ to the group $G_m^1 \times G$, $G_m^1 = GL(m)$. According to our recent result, [8], this reduction defines a principal torsion-free connection $\mathcal{E}_r(\Gamma, \Lambda)$ on $W^r P$. In the special case G is a singleton, so that $W^r P$ is the r -th order frame bundle $P^r M$ of M , we reobtain the exponential operator from [7].

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Further we clarify that every reduction of $W^r P$ to $G_m^1 \times G$ identifies the adjoint bundle of $W^r P$ with a vector bundle $\mathcal{L}^r P$ associated to $P^1 M \times_M P$. Then we outline how these results can be applied to the problem of finding all naturally induced connection on $W^r P$. Finally, we deduce that $\exp^{\Gamma, \Lambda}$ can be interpreted as the exponential map of the classical connection $N(\Gamma, \Lambda)$ on P constructed from Γ and Λ in [6].

All manifolds and maps are assumed to be infinitely differentiable. Unless otherwise specified, we use the terminology and notation from the book [9].

1. The exponential map.

By a classical connection Λ on a manifold M we mean a linear connection on TM . In the jet form, Λ is a linear morphism

$$\Lambda: TM \rightarrow J^1 TM.$$

Consider a principal G -bundle $p: P \rightarrow M$, a principal connection Γ on P and a classical torsion-free connection Λ on M . So Λ defines a local map \exp_x^Λ of $T_x M$ into M for every $x \in M$. In what follows we always assume that the vectors in question lie in the domain of definition of the exponential map under consideration. For a vector $X \in T_x M$, we construct the Γ -lift $\gamma(t)$ at $u \in P$ of the curve $\exp_x^\Lambda(tX)$ on M , $x = p(u)$, and we define

$$\lambda(\Gamma, \Lambda)(X, u) = \gamma(1) \in P.$$

For a vector $Z \in T_u P$, we consider $Z_0 = Tp(Z) \in T_x M$ and $\omega Z \in \mathfrak{g}$, where ω is the connection form of Γ . The exponential map \exp^G of G maps ωZ into an element of G and we define

$$\exp_u^{\Gamma, \Lambda}(Z) = \lambda(\Gamma, \Lambda)(Z_0, u) \exp^G(\omega Z).$$

This yields a local map $\exp_u^{\Gamma, \Lambda}$ of $T_u P$ into P . Clearly, $p \circ \exp_u^{\Gamma, \Lambda} = \exp_x^\Lambda \circ Tp$.

Lemma 1. *For every $u \in P$ and $g \in G$, the following diagram commutes*

$$(1) \quad \begin{array}{ccc} T_u P & \xrightarrow{\exp_u^{\Gamma, \Lambda}} & P \\ Tr_g \downarrow & & \downarrow r_g \\ T_{ug} P & \xrightarrow{\exp_{ug}^{\Gamma, \Lambda}} & P \end{array}$$

Proof. We have $\lambda(\Gamma, \Lambda)(X, ug) = \lambda(\Gamma, \Lambda)(X, u)g$ and $\omega(Tr_g(Z)) = \text{Ad}(g^{-1})(\omega Z)$. For every $A \in \mathfrak{g}$, $\exp^G(\text{Ad}(g^{-1})(A)) = g^{-1} \exp^G(A)g$. Hence $\exp_{ug}^{\Gamma, \Lambda}(Tr_g(Z)) = \gamma(1)g \exp^G(\text{Ad}(g^{-1})(\omega Z)) = \gamma(1) \exp^G(\omega Z)g$. \square

Every frame $v \in P_x^1 M$ is a linear map $v: \mathbb{R}^m \rightarrow T_x M$. We introduce a local map $\exp_{v,u}^{\Gamma,\Lambda}$ of $\mathbb{R}^m \times G$ into P by

$$(2) \quad \exp_{v,u}^{\Gamma,\Lambda}(y, g) = \lambda(\Gamma, \Lambda)(v(y), u)g, \quad y \in \mathbb{R}^m, \quad g \in G.$$

Clearly, this is a local \mathcal{PB} -isomorphism and we have

$$(3) \quad \exp_{vh,ug}^{\Gamma,\Lambda} = \exp_{v,u}^{\Gamma,\Lambda} \circ (l(h) \times l_g), \quad h \in G_m^1, \quad g \in G,$$

where $l(h): \mathbb{R}^m \rightarrow \mathbb{R}^m$ is the linear map determined by $h \in G_m^1$ and $l_g: G \rightarrow G$ is the left translation.

2. The exponential operator.

Hence the rule

$$(4) \quad E_r(\Gamma, \Lambda)(v, u) = j_{(0,e)}^{r+1} \exp_{v,u}^{\Gamma,\Lambda}$$

defines a map $E_r(\Gamma, \Lambda): P^1 M \times_M P \rightarrow W^{r+1} P$. The structure group of $W^{r+1} P$ is $W_m^{r+1} G = G_m^{r+1} \times T_m^{r+1} G = W_0^{r+1}(\mathbb{R}^m \times G)$, [9], and we have a canonical injection

$$(5) \quad i_m^{r+1}: G_m^1 \times G \rightarrow W_m^{r+1} G, \quad i_m^{r+1}(h, g) = j_{(0,e)}^{r+1}(l(h) \times l_g).$$

Proposition 1. $E_r(\Gamma, \Lambda)(P^1 M \times_M P)$ is a reduction of $W^{r+1} P$ to $i_m^{r+1}(G_m^1 \times G)$.

Proof. This follows directly from (3) and (5). \square

By Proposition 5.2 of [8], $E_r(\Gamma, \Lambda)$ corresponds to a torsion-free connection on $W^r P$, that will be denoted by $\mathcal{E}_r(\Gamma, \Lambda)$. (We recall that the torsion of a connection on $W^r P$ is defined as the covariant exterior differential of the canonical one-form $TW^r P \rightarrow \mathbb{R}^m \times \mathfrak{w}_m^{r-1} G$.) The natural operator \mathcal{E}_r transforming (Γ, Λ) into a principal connection on $W^r P$ will be called the exponential operator. In the case $G = \{pt\}$, we obtain the exponential operator of $P^r M$, [7].

3. The adjoint bundle of $W^r P$.

We write

$$(6) \quad \mathcal{S}^r \mathbb{R}^{m*} = \bigoplus_{k=1}^r \mathcal{S}^k \mathbb{R}^{m*}.$$

For every vector space V , the (m, r) -velocities bundle $T_m^r V$ is decomposed, by means of translations, into the form $T_m^r V = V \times (T_m^r V)_0$ and a classical result reads

$$(7) \quad (T_m^r V)_0 = V \otimes \mathcal{S}^r \mathbb{R}^{m*}.$$

In particular, $(T_m^r \mathbb{R})_0 = \mathcal{S}^r \mathbb{R}^{m^*}$. The group G_m^r acts on the right on $(T_m^r V)_0$ by the jet composition. The “linear” injection $l_r: G_m^1 \rightarrow G_m^r$ is defined by

$$l_r(h) = j_0^r l(h).$$

Another classical result reads that the induced left action of G_m^1 on $\mathcal{S}^r \mathbb{R}^{m^*}$ is the standard tensorial one; we denote it by $\tau_r(h): \mathcal{S}^r \mathbb{R}^{m^*} \rightarrow \mathcal{S}^r \mathbb{R}^{m^*}$. In the case of (7), we obtain the tensor product $\text{id}_V \otimes \tau_r(h)$.

The canonical injection $\nu_m^r: G \rightarrow T_m^r G$ is of the form

$$\nu_m^r(g) = j_0^r \hat{g},$$

where \hat{g} is the constant map of \mathbb{R}^m into $g \in G$. The Lie algebra $\mathfrak{t}_m^r G$ of $T_m^r G$ coincides with $T_m^r \mathfrak{g}$. One verifies directly that $\text{Ad}(j_0^r \hat{g}): T_m^r \mathfrak{g} \rightarrow T_m^r \mathfrak{g}$ is of the form

$$(8) \quad \text{Ad}(g) \times \text{Ad}(g) \otimes \text{id}_{\mathcal{S}^r \mathbb{R}^{m^*}}: \mathfrak{g} \times (T_m^r \mathfrak{g})_0 \rightarrow \mathfrak{g} \times (T_m^r \mathfrak{g})_0.$$

In the decomposition $W_m^r G = G_m^r \times T_m^r G$, the injection $i_m^r: G_m^1 \times G \rightarrow W_m^r G$ is of the form

$$i_m^r(h, g) = (l_r(h), \nu_m^r(g)).$$

As a vector space, the Lie algebra of $W_m^r G$ is the product

$$\mathfrak{w}_m^r G = \mathfrak{g}_m^r \times \mathfrak{t}_m^r G.$$

But $\mathfrak{g}_m^r \approx (T_m^r \mathbb{R}^m)_0 = \mathbb{R}^m \otimes \mathcal{S}^r \mathbb{R}^{m^*}$ and $\mathfrak{t}_m^r G = T_m^r \mathfrak{g} = \mathfrak{g} \times \mathfrak{g} \otimes \mathcal{S}^r \mathbb{R}^{m^*}$, so that we have

$$(9) \quad \mathfrak{w}_m^r G = \mathbb{R}^m \otimes \mathcal{S}^r \mathbb{R}^{m^*} \times \mathfrak{g} \times \mathfrak{g} \otimes \mathcal{S}^r \mathbb{R}^{m^*}.$$

The proof of the following assertion is straightforward.

Lemma 2. *For every $(A_1, A_2, A_3) \in \mathbb{R}^m \otimes \mathcal{S}^r \mathbb{R}^{m^*} \times \mathfrak{g} \times \mathfrak{g} \otimes \mathcal{S}^r \mathbb{R}^{m^*}$, we have*

$$\begin{aligned} \text{Ad}(j_0^r l(h), j_0^r \hat{g})(A_1, A_2, A_3) = \\ ((l(h) \otimes \tau_r(h))(A_1), \text{Ad}(g)(A_2), (\text{Ad}(g) \otimes \tau_r(h))(A_3)). \end{aligned}$$

Analogously to (6), we write

$$\mathcal{S}^r T^* M = T^* M \times_M \cdots \times_M \mathcal{S}^r T^* M.$$

Let $\text{Ad } P = P[\mathfrak{g}, \text{Ad}]$ be the adjoint bundle of P . We introduce

$$(10) \quad \mathcal{L}^r P = TM \otimes \mathcal{S}^r T^* M \times_M \text{Ad } P \times_M \text{Ad } P \otimes \mathcal{S}^r T^* M.$$

Clearly, $\mathcal{L}^r P$ is a vector bundle associated to $P^1 M \times_M P$.

Consider an arbitrary reduction $I_r: P^1 M \times_M P \rightarrow W^r P$ to the subgroup $i_m^r(G_m^1 \times G)$. Then the results of this section can be summarized by

Proposition 2. *If we consider $\text{Ad}(W^r P)$ as a vector bundle associated to $I_r(P^1 M \times_M P)$, then $\text{Ad}(W^r P)$ coincides with $\mathcal{L}^r P$.*

4. Application to natural operators.

The difference of two principal connections on $W^r P$ is an arbitrary section of $\text{Ad}(W^r P) \otimes T^* M$, [9]. Beside the natural connection $\mathcal{E}_r(\Gamma, \Lambda)$, our construction yields a natural reduction $E_{r-1}(\Gamma, \Lambda)(P^1 M \times_M P)$ of $W^r P$ to the group $i_m^r(G_m^1 \times G)$, that identifies $\text{Ad} W^r P$ with $\mathcal{L}^r P$. This implies

Proposition 3. *The natural operators transforming (Γ, Λ) into principal connections on $W^r P$ are in bijection with the natural operators transforming (Γ, Λ) into sections of $\mathcal{L}^r P \otimes T^* M$.*

The latter operators can be determined by means of a generalization of the Utiyama theorem due to J. Janyška, [4]. Using [4] and our construction of the exponential operator, J. Janyška and J. Vondra, [5], found all natural operators transforming (Γ, Λ) into principal connections on $W^r P$. We remark that some further aspects of the same problem are discussed in [2].

5. Another construction of $\exp^{\Gamma, \Lambda}$.

We find remarkable that $\exp^{\Gamma, \Lambda}$ coincides with the exponential map of the classical connection $N(\Gamma, \Lambda)$ on P constructed from Γ and Λ in [6], see also Section 54.7 of [9]. We start with a slightly more general form of that construction.

Consider an arbitrary fibered manifold $p: Y \rightarrow M$, a general connection Γ on Y , a classical torsion-free connection Λ on M , a vector space W , $\dim W = \dim Y - \dim M$, and a vertical parallelism $\Phi: Y \times W \rightarrow VY$, that is an \mathcal{FM} -morphism over Y restricting to a linear isomorphism $W \rightarrow V_y Y$ for every $y \in Y$. So every $w \in W$ defines a vertical vector field $\varphi(w)$ on Y . In this situation, we construct a classical connection $(\Gamma, \Lambda, \Phi): TY \rightarrow J^1(TY \rightarrow Y)$ as follows. We decompose $X \in T_y Y$ into $hX = \Gamma(y, X_0)$, $X_0 \in T_x M$, $x = p(y)$, and $vX = \Phi(y, X_1)$, $X_1 \in W$. We take a vector field ξ on M such that $j_x^1 \xi = \Lambda(X_0)$ and consider its Γ -lift $\Gamma\xi: Y \rightarrow TY$. Then we define

$$(11) \quad (\Gamma, \Lambda, \Phi)(X) = j_y^1(\Gamma\xi + \varphi(X_1)).$$

In fiber coordinates (x^i, y^p) on Y , we have

$$(12) \quad \Gamma \equiv dy^p = F_i^p(x, y) dx^i,$$

$$(13) \quad \Lambda \equiv d\xi^i = \Lambda_{kj}^i(x) \xi^k dx^j, \quad \xi^i = dx^i,$$

$$(14) \quad \Phi \equiv \eta^p = a_q^p(x, y) w^q, \quad \eta^p = dy^p, \quad (w^q) \in W.$$

Consider $(x_0, y_0, \xi_0, \eta_0) \in TY$. If $\xi \equiv (\xi^i(x))$, then $\Gamma\xi \equiv (\xi^i(x), F_i^p(x, y)\xi^i)$ and $\varphi(X_1) = a_r^p(x, y)\tilde{a}_q^r(x_0, y_0)(\eta_0^q - F_i^q(x_0, y_0)\xi_0^i)$, where \tilde{a}_q^p means the inverse matrix to a_q^p . Passing to 1-jets, we obtain the equations of (Γ, Λ, Φ) in the form (13) and

$$(15) \quad \begin{aligned} d\eta^p &= \left(\frac{\partial F_i^p}{\partial x^j} + F_k^p \Lambda_{ij}^k \right) \xi^i dx^j + \frac{\partial F_i^p}{\partial y^q} \xi^i dy^q \\ &+ \frac{\partial a_r^p}{\partial x^j} \tilde{a}_q^r (\eta^q - F_i^q \xi^i) dx^j + \frac{\partial a_s^p}{\partial y^q} \tilde{a}_r^s (\eta^r - F_i^r \xi^i) dy^q. \end{aligned}$$

Lemma 3. *Every Γ -lift $(x(t), y(t))$ of a geodesic $x(t)$ of Λ is a geodesic of (Γ, Λ, Φ) for arbitrary Φ .*

Proof. The Γ -lift satisfies

$$(16) \quad \frac{dy^p}{dt} = F_i^p(x(t), y(t)) \frac{dx^i}{dt}.$$

Differentiating (16) and using $x(t)$ is a geodesic of Λ , we find

$$(17) \quad \frac{d^2 y^p}{dt^2} = \left(\frac{\partial F_i^p}{\partial x^j} + F_k^p \Lambda_{ij}^k \right) \frac{dx^i}{dt} \frac{dx^j}{dt} + \frac{\partial F_i^p}{\partial y^q} \frac{dx^i}{dt} \frac{dy^q}{dt}.$$

But (16) and (17) annihilates the equations of geodesics corresponding to (15) for every Φ . \square

If P is a principal bundle, Γ is a principal connection on P , $W = \mathfrak{g}$ and Φ is the canonical vertical parallelism of P , then (Γ, Λ, Φ) coincides with $N(\Gamma, \Lambda)$ from [6]. Then our claim that $\exp^{\Gamma, \Lambda} = \exp^{N(\Gamma, \Lambda)}$ follows from the following assertion.

Proposition 4. *If $z(t)$ is a Γ -lift of a geodesic $x(t)$ of Λ and $g(t)$ is a one-parameter subgroup of G , then $z(t)g(t)$ is also a geodesic of $N(\Gamma, \Lambda)$.*

Proof. By locality and by the Ado theorem, it suffices to discuss the case $P = \mathbb{R}^m \times GL(n)$. Then we have the coordinate expressions

$$(18) \quad \Gamma \equiv dy_q^p = \Gamma_{ri}^p(x) y_q^r dx^i,$$

$$(19) \quad \Phi \equiv \eta_q^p = y_r^p w_q^r.$$

The one-parameter subgroups of $GL(n)$ are e^{At} with arbitrary $A \in \mathfrak{gl}(n)$. The Γ -lift $z(t) = (x^i(t), y_q^p(t))$ of $x(t) = (x^i(t))$ satisfies (18). In this situation, one verifies that $z(t)e^{At}$ is also a geodesic by direct, but tedious evaluations. \square

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